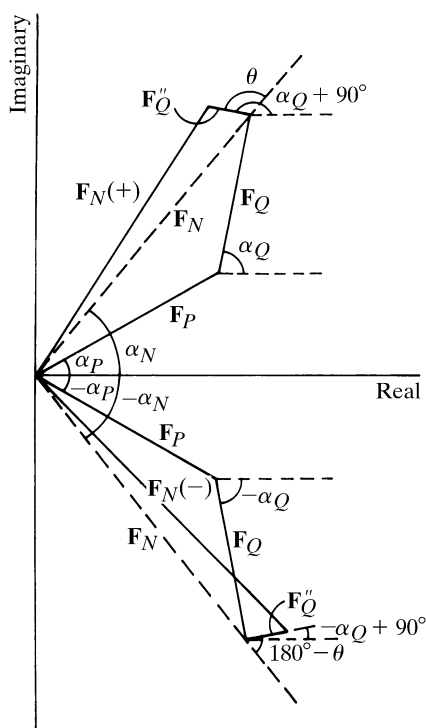


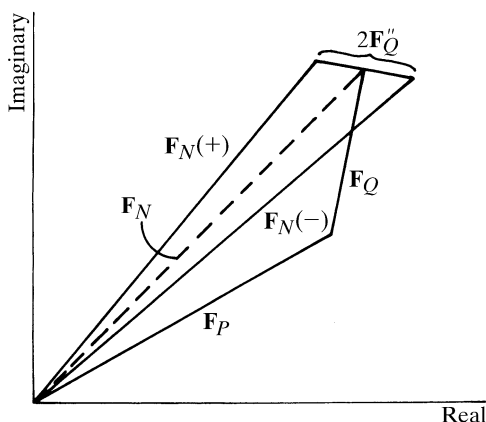
## 2.4. ISOMORPHOUS REPLACEMENT AND ANOMALOUS SCATTERING


 Fig. 2.4.3.2. Vector diagram illustrating the violation of Friedel's law when  $F_Q'' \neq 0$ .

does not give rise to intensity differences. Expressions for intensity differences between Friedel equivalents have been derived by Zachariasen (1965) for the most general case of a structure containing normal as well as different types of anomalous scatterers. Statistical distributions of such differences under various conditions have also been derived (Parthasarathy & Srinivasan, 1964; Parthasarathy, 1967). It turns out that, with a single type of anomalous scatterer in the structure, the ratio

$$\frac{|F_N^2(+)-F_N^2(-)|}{F_N^2(+)+F_N^2(-)}$$

has a maximum mean value when the scattering powers of the anomalous scatterers and the normal scatterers are nearly the same (Srinivasan, 1972). Also, for a given ratio between the scattering powers, the smaller the number of anomalous scatterers, the higher is the mean ratio.


 Fig. 2.4.3.3. A composite view of the vector relationship between  $F_N(+)$  and  $F_N(-)$ .

## 2.4.3.3. Friedel and Bijvoet pairs

The discussion so far has been concerned essentially with crystals belonging to space groups  $P1$  and  $P\bar{1}$ . In the centrosymmetric space group, the crystal and the diffraction pattern have the same symmetry, namely, an inversion centre. In  $P1$ , however, the crystals are noncentrosymmetric while the diffraction pattern has an inversion centre, in the absence of anomalous scattering. When anomalous scatterers are present in the structure ( $F_Q'' \neq 0$ ), Friedel's law breaks down and the diffraction pattern no longer has an inversion centre. Thus the diffraction pattern displays the same symmetry as that of the crystal in the presence of anomalous scattering. The same is true with higher-symmetry space groups also. For example, consider a crystal with space group  $P222$ , containing anomalous scatterers. The magnitudes of  $F_P$  are the same for all equivalent reflections; so are those of  $F_Q$  and  $F_Q''$ . Their phase angles, however, differ from one equivalent to another, as can be seen from Table 2.4.3.1. When  $F_Q'' = 0$ , the magnitudes of the vector sum of  $F_P$  and  $F_Q$  are the same for all the equivalent reflections. The intensity pattern thus has point-group symmetry  $2/m\ 2/m\ 2/m$ . When  $F_Q'' \neq 0$ , the equivalent reflections can be grouped into two sets in terms of their intensities:  $hkl, h\bar{k}\bar{l}, \bar{h}k\bar{l}$  and  $\bar{h}\bar{k}l$ ; and  $\bar{h}\bar{k}l, \bar{h}kl, h\bar{k}l$  and  $hkl$ . The equivalents belonging to the first group have the same intensity; so have the equivalents in the second group. But the two intensities are different. Thus the symmetry of the pattern is  $222$ , the same as that of the crystal.

In general, under conditions of anomalous scattering, equivalent reflections generated by the symmetry elements in the crystal have intensities different from those of equivalent reflections generated by the introduction of an additional inversion centre in normal scattering. There have been suggestions that a reflection from the first group and another from the second group should be referred to as a 'Bijvoet pair' instead of a 'Friedel pair', when the two reflections are not inversely related. Most often, however, the terms are used synonymously. The same practice will be followed in this article.

## 2.4.3.4. Determination of absolute configuration

The determination of the absolute configuration of chiral molecules has been among the most important applications of anomalous scattering. Indeed, anomalous scattering is the only effective method for this purpose and the method, first used in the early 1950s (Peerdeman *et al.*, 1951), has been extensively employed in structural crystallography (Ramaseshan, 1963; Vos, 1975).

Many molecules, particularly biologically important ones, are chiral in that the molecular structure is not superimposable on its mirror image. Chirality (handedness) arises primarily on account of the presence of asymmetric carbon atoms in the molecule. A tetravalent carbon is asymmetric when the four atoms (or groups) bonded to it are all different from one another. The substituents can then have two distinct arrangements which are mirror images of (or related by inversion to) each other. These optical isomers or enantiomers have the same chemical and physical properties except

 Table 2.4.3.1. Phase angles of different components of the structure factor in space group  $P222$ 

Reflection	Phase angle ( $^\circ$ ) of		
	$F_P$	$F_Q$	$F_Q''$
$hkl, h\bar{k}\bar{l}, \bar{h}k\bar{l}, \bar{h}\bar{k}l$	$\alpha_P$	$\alpha_Q$	$90 + \alpha_Q$
$\bar{h}\bar{k}l, \bar{h}kl, h\bar{k}l, hkl$	$-\alpha_P$	$-\alpha_Q$	$90 - \alpha_Q$