

## 4.6. RECIPROCAL-SPACE IMAGES OF APERIODIC CRYSTALS

$$(S_m^{-1})^T = \frac{1}{2+\tau} \begin{pmatrix} \tau^2 - x\tau + 1 & -x\tau^2 \\ x & \tau^2 + x\tau + 1 \end{pmatrix}_D,$$

where  $x = (n\tau - m)/(m\tau + n)$ :

$$\begin{aligned} \mathbf{d}_i^* &= \sum_{j=1}^2 (S_m^{-1})_{ij}^T \mathbf{d}_j^*; \\ \mathbf{d}_1^* &= \frac{1}{2+\tau} \left[ (\tau^2 - x\tau + 1) \mathbf{d}_1^* - x\tau^2 \mathbf{d}_2^* \right] \\ &= a^* \begin{pmatrix} 1 - x\tau \\ -\tau \end{pmatrix}_V \\ &= a^* \begin{pmatrix} 2m\tau - n\tau \\ m\tau + n \\ -\tau \end{pmatrix}_V, \\ \mathbf{d}_2^* &= \frac{1}{2+\tau} \left[ x\mathbf{d}_1^* + (\tau^2 + x\tau + 1) \mathbf{d}_2^* \right] \\ &= a^* \begin{pmatrix} \tau + x \\ 1 \end{pmatrix}_V \\ &= a^* \begin{pmatrix} 2n\tau + m\tau \\ m\tau + n \\ 1 \end{pmatrix}_V. \end{aligned}$$

The point  $x_n(t)$  of the  $n$ th interval L or S of an infinite Fibonacci sequence is given by

$$x_n(t) = \{x_0 + n(3 - \tau) - (\tau - 1)[\text{frac}(n\tau + t) - (1/2)]\}S,$$

where  $t$  is the phase of the modulation function  $y(t) = (\tau - 1)[\text{frac}(n\tau + t) - (1/2)]$  (Janssen, 1986). Thus, the Fibonacci sequence can also be dealt with as an incommensurately modulated structure. This is a consequence of the fact that for 1D structures only the crystallographic point symmetries 1 and  $\bar{1}$  allow the existence of a periodic average structure.

The embedding of the Fibonacci chain as an incommensurately modulated structure can be performed as follows:

(1) select a subset  $\Lambda^* \subset M^*$  of strong reflections for main reflections  $\mathbf{H} = h\mathbf{a}^*, h \in \mathbb{Z}$ ;

(2) define a satellite vector  $\mathbf{q} = \alpha\mathbf{a}^*$  pointing from each main reflection to the next satellite reflection.

One possible way of indexing based on the same  $\mathbf{a}^*$  as defined above is illustrated in Fig. 4.6.2.10. The scattering vector is given by  $\mathbf{H}^\parallel = h(\tau + 1)\mathbf{a}^* + mq$ , where  $\mathbf{q} = \tau\mathbf{a}^*$ , or, in the 2D representation,  $\mathbf{H} = h_1\mathbf{d}_1^* + h_2\mathbf{d}_2^*$ , where  $\mathbf{d}_1^* = a^* \begin{pmatrix} 1 + \tau \\ 0 \end{pmatrix}_V$  and  $\mathbf{d}_2^* = a^* \begin{pmatrix} \tau \\ 1 \end{pmatrix}_V$ , with the direct basis

$$\mathbf{d}_1 = \frac{1}{a^*(1 + \tau)} \begin{pmatrix} 1 \\ -\tau \end{pmatrix}_V, \quad \mathbf{d}_2 = \frac{1}{a^*} \begin{pmatrix} 0 \\ 1 \end{pmatrix}_V.$$

The modulation function is saw-tooth-like (Fig. 4.6.2.11).

## 4.6.2.5. 1D structures with fractal atomic surfaces

A 1D structure with a *fractal atomic surface* (Hausdorff dimension 0.9157...) can be derived from the Fibonacci sequence by squaring its substitution matrix  $S$ :

$$\begin{pmatrix} S \\ L \end{pmatrix} \rightarrow \begin{pmatrix} 1 & 1 \\ 1 & 2 \end{pmatrix} \begin{pmatrix} S \\ L \end{pmatrix} = \begin{pmatrix} S + L \\ S + 2L \end{pmatrix}$$

$$\text{with } S^2 = \begin{pmatrix} 1 & 1 \\ 1 & 2 \end{pmatrix},$$

corresponding to the substitution rule  $S \rightarrow SL, L \rightarrow LLS$  as well

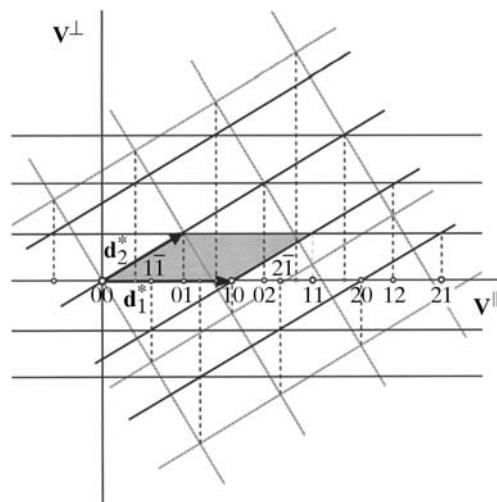


Fig. 4.6.2.10. Reciprocal space of the embedded Fibonacci chain as a modulated structure. Several main and satellite reflections are indexed. The square reciprocal lattice of the quasicrystal description illustrated in Fig. 4.6.2.9 is indicated by grey lines. The reflections located on  $\mathbf{V}^\parallel$  can be considered to be projected either from the 2D square lattice of the embedding as for a QS or from the 2D oblique lattice of the embedding as for an IMS.

as two other non-equivalent ones (see Janssen, 1995). The eigenvalues  $\lambda_i$  are obtained by calculating

$$\det |S - \lambda I| = 0.$$

The evaluation of the determinant gives the characteristic polynomial

$$\lambda^2 - 3\lambda + 1 = 0,$$

with the solutions  $\lambda_{1,2} = [3 \pm (5)^{1/2}]/2$ , with  $\lambda_1 = \tau^2$  and  $\lambda_2 = 1/\tau^2 = 2 - \tau$ , and the same eigenvectors  $\mathbf{w}_1 = \begin{pmatrix} 1 \\ \tau \end{pmatrix}, \mathbf{w}_2 = \begin{pmatrix} 1 \\ -1/\tau \end{pmatrix}$  as for the Fibonacci sequence. Rewriting the eigenvalue equation gives

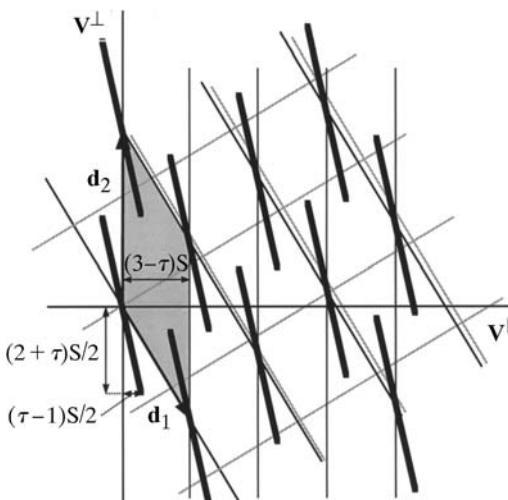


Fig. 4.6.2.11. 2D direct-space embedding of the Fibonacci chain as a modulated structure. The average period is  $(3 - \tau)S$ . The square lattice in the quasicrystal description shown in Fig. 4.6.2.8 is indicated by grey lines. The rod-like atomic surfaces are now inclined relative to  $\mathbf{V}^\parallel$  and arranged so as to give a saw-tooth modulation wave.