

## 1. GENERAL RELATIONSHIPS AND TECHNIQUES

Table 1.2.7.3. ‘Kubic Harmonic’ functions

(a) Coefficients in the expression  $K_{lj} = \sum_{mp} k_{mpj}^l y_{lmp}$  with normalization  $\int_0^\pi \int_0^{2\pi} |K_{lj}|^2 \sin \theta \, d\theta \, d\varphi = 1$  (Kara & Kurki-Suonio, 1981).

Even $l$		$mp$						
$l$	$j$	0+	2+	4+	6+	8+	10+	
0	1	1						
4	1	$\frac{1}{2} \left(\frac{7}{3}\right)^{1/2}$ 0.76376		$\frac{1}{2} \left(\frac{5}{3}\right)^{1/2}$ 0.64550				
6	1	$\frac{1}{2} \left(\frac{1}{2}\right)^{1/2}$ 0.35355		$-\frac{1}{2} \left(\frac{7}{2}\right)^{1/2}$ -0.93541				
6	2		$\frac{1}{4} 11^{1/2}$ 0.82916		$-\frac{1}{4} 5^{1/2}$ -0.55902			
8	1	$\frac{1}{8} 33^{1/2}$ 0.71807		$\frac{1}{4} \left(\frac{7}{3}\right)^{1/2}$ 0.38188		$\frac{1}{8} \left(\frac{65}{3}\right)^{1/2}$ 0.58184		
10	1	$\frac{1}{8} \left(\frac{65}{6}\right)^{1/2}$ 0.41143		$-\frac{1}{4} \left(\frac{11}{2}\right)^{1/2}$ -0.58630		$-\frac{1}{8} \left(\frac{187}{6}\right)^{1/2}$ -0.69784		
10	2		$\frac{1}{8} \left(\frac{247}{6}\right)^{1/2}$ 0.80202		$\frac{1}{16} \left(\frac{19}{3}\right)^{1/2}$ 0.15729		$\frac{1}{16} 85^{1/2}$ 0.57622	
$l$	$j$		2–	4–	6–	8–		
3	1		1					
7	1		$\frac{1}{2} \left(\frac{13}{6}\right)^{1/2}$ 0.73598		$\frac{1}{2} \left(\frac{11}{16}\right)^{1/2}$ 0.41458			
9	1		$\frac{1}{4} 3^{1/2}$ 0.43301		$-\frac{1}{4} 13^{1/2}$ -0.90139			
9	2		$\frac{1}{2} \left(\frac{17}{6}\right)^{1/2}$ 0.84163		$-\frac{1}{2} \left(\frac{7}{6}\right)^{1/2}$ -0.54006			

(b) Coefficients  $k_{mpj}^l$  and density normalization factors  $N_{lj}$  in the expression  $K_{lj} = N_{lj} \sum_{mp} k_{mpj}^l u_{lmp}$  where  $u_{lm\pm} = P_l^m(\cos \theta) \frac{\cos m\varphi}{\sin m\varphi}$  (Su & Coppens, 1994).

Even $l$		$N_{lj}$	$mp$						
$l$	$j$		0+	2+	4+	6+	8+	10+	
0	1	$1/4\pi = 0.079577$	1						
4	1	0.43454	1		$+1/168$				
6	1	0.25220	1		$-1/360$				
6	2	0.020833		1		$-1/792$			

## 1.2. THE STRUCTURE FACTOR

Table 1.2.7.3. ‘Kubic Harmonic’ functions (*cont.*)

Even $l$		$N_{lj}$	$mp$					
8	1	0.56292	1		1/5940			$\frac{1}{672} \times \frac{1}{5940}$
10	1	0.36490	1		1/5460			$\frac{1}{4320} \times \frac{1}{5460}$
10	2	0.0095165	1			1/43680		$-\frac{1}{456} \times \frac{1}{43680}$
$l$	$j$			2–	4–	6–	8–	
3	1	0.066667		1				
7	1	0.014612		1		1/1560		
9	1	0.0059569		1		1/2520		
9	2	0.00014800			1			-1/4080

(c) Density-normalized Kubic harmonics as linear combinations of density-normalized spherical harmonic functions. Coefficients in the expression  $K_{lj} = \sum_{mp} k''_{mpj} d_{lmp}$ . Density-type normalization is defined as  $\int_0^\pi \int_0^{2\pi} |K_{lj}| \sin \theta \, d\theta \, d\varphi = 2 - \delta_{l0}$ .

Even $l$		$mp$						
$l$	$j$	0+	2+	4+	6+	8+	10+	
0	1	1						
4	1	0.78245		0.57939				
6	1	0.37790		-0.91682				
6	2		0.83848		-0.50000			
$l$	$j$	2–	4–	6–	8–			
3	1	1						
7	1	0.73145		0.63290				

(d) Index rules for cubic symmetries (Kurki-Suonio, 1977; Kara & Kurki-Suonio, 1981).

$l$	$j$	23 $T$	$m\bar{3}$ $T_h$	432 $O$	$\bar{4}3m$ $T_d$	$m\bar{3}m$ $O_h$
0	1	×	×	×	×	×
3	1	×			×	
4	1	×	×	×	×	×
6	1	×	×	×	×	×
6	2	×	×			
7	1	×			×	
8	1	×	×	×	×	×
9	1	×			×	
9	2	×		×		
10	1	×	×	×	×	×
10	2	×	×			

### 1.2.8. Fourier transform of orbital products

by (Stewart, 1969*a*)

If the wavefunction is written as a sum over normalized Slater determinants, each representing an antisymmetrized combination of occupied molecular orbitals  $\chi_i$  expressed as linear combinations of atomic orbitals  $\varphi_\nu$ , i.e.  $\chi_i = \sum_\nu c_{i\nu} \varphi_\nu$ , the electron density is given

$$\rho(\mathbf{r}) = \sum_i n_i \chi_i^2 = \sum_\mu \sum_\nu P_{\mu\nu} \varphi_\mu(\mathbf{r}) \varphi_\nu(\mathbf{r}), \quad (1.2.8.1)$$

with  $n_i = 1$  or 2. The coefficients  $P_{\mu\nu}$  are the populations of the