International Tables for Crystallography (2006). Vol. D, Section 1.2.2.2, pp. 36-37.

1. TENSORIAL ASPECTS OF PHYSICAL PROPERTIES

(c) point groups that have elements in common with $O(3)\SO(3)$ but do not contain -E; such groups are isomorphic to a group of the first class, as one can see if one multiplies all elements with determinant equal to -1 by -E.

The list of three-dimensional finite point groups is given in Table 1.2.6.1. All isomorphism classes of two-dimensional point groups occur in three dimensions as well. The isomorphism classes occurring here for the first time are:

 $C_n \times C_2: A, B, \text{ with } A^n = B^2 = ABA^{-1}B^{-1} = E;$ $D_n \times C_2: A, B, C \text{ with } A^n = B^2 = (AB)^2 = C^2 = ACA^{-1}C^{-1}$ $= BCB^{-1}C^{-1} = E;$ $T: A, B, \text{ with } A^3 = B^2 = (AB)^3 = E;$ $O: A, B, \text{ with } A^4 = B^3 = (AB)^2 = E;$ $T \times C_2: A, B, C, \text{ with } A^3 = B^2 = (AB)^3 = C^2 = ACA^{-1}C^{-1}$ $= BCB^{-1}C^{-1} = E;$ $O \times C_2: A, B, C, \text{ with } A^4 = B^3 = (AB)^2 = C^2 = ACA^{-1}C^{-1}$ $= BCB^{-1}C^{-1} = E;$ $I: A, B, \text{ with } A^5 = B^3 = (AB)^2 = E;$ $I \times C_2: A, B, C, \text{ with } A^5 = B^3 = (AB)^2 = C^2 = ACA^{-1}C^{-1}$ $= BCB^{-1}C^{-1} = E.$

The crystallographic groups among them are given in Table 1.2.6.2.

1.2.2.2. Representations of finite groups

As stated in Section 1.2.1, elements of point groups act on physical properties (like tensorial properties) and on wave functions as linear operators. These linear operators therefore generally act in a different space than the three-dimensional configuration space. We denote this new space by V and consider a mapping D from the point group K to the group of nonsingular linear operators in V that satisfies

$$D(R)D(R') = D(RR') \quad \forall R, R' \in K.$$
 (1.2.2.3)

In other words D is a homomorphism from K to the group of nonsingular linear transformations GL(V) on the vector space V. Such a homomorphism is called a *representation* of K in V. Here we only consider finite-dimensional representations.

With respect to a basis \mathbf{e}_i (i = 1, 2, ..., n) the linear transformations are given by matrices $\Gamma(R)$. The mapping Γ from K to the group of nonsingular $n \times n$ matrices GL(n, R) (for a real vector space V) or GL(n, C) (if V is complex) is called an *n*-dimensional *matrix representation* of K.

If one chooses another basis for V connected to the former one by a nonsingular matrix S, the same group of operators D(K) is represented by another matrix group $\Gamma'(K)$, which is related to $\Gamma(K)$ by S according to $\Gamma'(R) = S^{-1}\Gamma(R)S$ ($\forall R \in K$). Two such matrix representations are called *equivalent*. On the other hand, two such equivalent matrix representations can be considered to describe two different groups of linear operators [D(K) and D'(K)] on the same basis. Then there is a nonsingular linear operator T such that D(R)T = TD'(R) ($\forall R \in K$). In this case, the representations D(K) and D'(K) are also called equivalent.

It may happen that a representation D(K) in V leaves a subspace W of V invariant. This means that for every vector $v \in W$ and every element $R \in K$ one has $D(R)v \in W$. Suppose that this subspace is of dimension m < n. Then one can choose m basis vectors for V inside the invariant subspace. With respect to this basis, the corresponding matrix representation has elements

$$\Gamma(R) = \begin{pmatrix} \Gamma_1(R) & \Gamma_3(R) \\ 0 & \Gamma_2(R) \end{pmatrix}, \qquad (1.2.2.4)$$

where the matrices $\Gamma_1(R)$ form an *m*-dimensional matrix representation of *K*. In this situation, the representations D(K) and $\Gamma(K)$ are called *reducible*. If there is no proper invariant subspace the representation is *irreducible*. If the representation is a direct sum of subspaces, each carrying an irreducible representation, the representation is called *fully reducible* or *decomposable*. In the latter case, a basis in *V* can be chosen such that the matrices $\Gamma(R)$ are direct sums of matrices $\Gamma_i(R)$ such that the $\Gamma_i(R)$ form an irreducible matrix representation. If $\Gamma_3(R)$ in (1.2.2.4) is zero and Γ_1 and Γ_2 form irreducible matrix representations, Γ is fully reducible. For finite groups, each reducible representation is fully reducible. That means that if $\Gamma(K)$ is reducible, there is a matrix *S* such that

$$\Gamma(R) = S \begin{bmatrix} \Gamma_1(R) \oplus \dots \oplus \Gamma_n(R) \end{bmatrix} S^{-1}$$

$$= S \begin{pmatrix} \Gamma_1(R) & 0 & \dots & 0 \\ 0 & \Gamma_2(R) & \dots & 0 \\ \vdots & \vdots & \ddots & \vdots \\ 0 & 0 & \dots & \Gamma_n(R) \end{pmatrix} S^{-1}.$$
(1.2.2.5)

In this way one may proceed until all matrix representations $\Gamma_i(K)$ are *irreducible*, *i.e.* do not have invariant subspaces. Then each representation $\Gamma(K)$ can be written as a direct sum

$$\Gamma(R) = S[m_1 \Gamma_1(R) \oplus \ldots \oplus m_s \Gamma_s(R)]S^{-1}, \qquad (1.2.2.6)$$

where the representations $\Gamma_1 \dots \Gamma_s$ are all nonequivalent and the *multiplicities* m_i are the numbers of times each irreducible representation occurs. The nonequivalent irreducible representations Γ_i for which the multiplicity is not zero are the *irreducible components* of $\Gamma(K)$.

We first discuss two special representations. The simplest representation in one-dimensional space is obtained by assigning the number 1 to all elements of K. Obviously this is a representation, called the *identity* or *trivial representation*. Another is the *regular representation*. To obtain this, one numbers the elements of K from 1 to the order N of the group (|K| = N). For a given $R \in K$ there is a one-to-one mapping from K to itself defined by $R_i \rightarrow R_j \equiv RR_i$. Consider the $N \times N$ matrix $\Gamma(R)$, which has in the *i*th column zeros except on line *j*, where the entry is unity. The matrix $\Gamma(R)$ then has as only entries 0 or 1 and satisfies

$$RR_i = \Gamma(R)_{ii}R_i, \quad (i = 1, 2, \dots, N).$$
(1.2.2.7)

These matrices $\Gamma(R)$ form a representation, the *regular representation* of K of dimension N, as one sees from

$$(R_{i}R_{j})R_{k} = R_{i}\sum_{l=1}^{N}\Gamma(R_{j})_{lk}R_{l} = \sum_{l=1}^{N}\sum_{m=1}^{N}\Gamma(R_{j})_{lk}\Gamma(R_{i})_{ml}R_{m}$$
$$= \sum_{m=1}^{N}[\Gamma(R_{i})\Gamma(R_{j})]_{mk}R_{m} = \sum_{m=1}^{N}\Gamma(R_{i}R_{j})_{mk}R_{m}.$$

A representation in a real vector space that leaves a positive definite metric invariant can be considered on an orthonormal basis for that metric. Then the matrices satisfy

 $\Gamma(R)\Gamma(R)^T = E$

 $(^{T}$ denotes transposition of the matrix) and the representation is *orthogonal*. If V is a complex vector space with positive definite metric invariant under the representation, the latter gives on an orthonormal basis matrices satisfying

$$\Gamma(R)\Gamma(R)^{\dagger} = E$$

1.2. REPRESENTATIONS OF CRYSTALLOGRAPHIC GROUPS

([†] denotes Hermitian conjugation) and the representation is *unitary*. A real representation of a finite group is always equivalent with an orthogonal one, a complex representation of a finite group is always equivalent with a unitary one. As a proof of the latter statement, consider the standard Hermitian metric on $V: f(x, y) = \sum_i x_i^* y_i$. Then the positive definite form

$$F(x, y) = (1/N) \sum_{R \in K} f(D(R)x, D(R)y)$$
(1.2.2.8)

is invariant under the representation. To show this, take an arbitrary element R'. Then

$$F(D(R')x, D(R')y) = (1/N) \sum_{R \in K} f(D(R'R)x, D(R'R)y)$$

= F(x, y). (1.2.2.9)

With respect to an orthonormal basis for this metric F(x, y), the matrices corresponding to D(R) are unitary. The complex representation can be put into this unitary form by a basis transformation. For a real representation, the argument is fully analogous, and one obtains an orthogonal transformation.

From two representations, $D_1(K)$ in V_1 and $D_2(K)$ in V_2 , one can construct the sum and product representations. The *sum representation* acts in the direct sum space $V_1 \oplus V_2$, which has elements (\mathbf{a}, \mathbf{b}) with $\mathbf{a} \in V_1$ and $\mathbf{b} \in V_2$. The representation $D_1 \oplus D_2$ is defined by

$$[(D_1 \oplus D_2)(R)](\mathbf{a}, \mathbf{b}) = (D_1(R)\mathbf{a}, D_2(R)\mathbf{b}).$$
(1.2.2.10)

The matrices $\Gamma_1 \oplus \Gamma_2(R)$ are of dimension $n_1 + n_2$.

The *product representation* acts in the tensor space, which is the space spanned by the vectors $\mathbf{e}_i \otimes \mathbf{e}_j$ $(i = 1, 2, ..., \dim V_1; j = 1, 2, ..., \dim V_2)$. The dimension of the tensor space is the product of the dimensions of both spaces. The action is given by

$$[(D_1 \otimes D_2)(R)] \mathbf{a} \otimes \mathbf{b} = D_1(R) \mathbf{a} \otimes D_2(R) \mathbf{b}.$$
(1.2.2.11)

For bases \mathbf{e}_i $(i = 1, 2, ..., d_1)$ for V_1 and \mathbf{e}'_j $(j = 1, 2, ..., d_2)$ for V_2 , a basis for the tensor product of spaces is given by

$$\mathbf{e}_i \otimes \mathbf{e}'_j, \quad i = 1, \dots, d_1; \ j = 1, 2, \dots, d_2,$$
 (1.2.2.12)

and with respect to this basis the representation of K is given by matrices

$$(\Gamma_1 \otimes \Gamma_2)(R)_{ik,jl} = \Gamma_1(R)_{ij}\Gamma_2(R)_{kl}.$$
 (1.2.2.13)

As an example of these operations, consider

$$\begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \oplus \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & 1 & 0 \end{pmatrix};$$
$$\begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \otimes \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} = \begin{pmatrix} 0 & 1 & 0 & 0 \\ 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & -1 \\ 0 & 0 & -1 & 0 \end{pmatrix}.$$

If two representations $D_1(K)$ and $D_2(K)$ are equivalent, there is an operator S such that

$$SD_1(R) = D_2(R)S \quad \forall R \in K.$$

This relation may also hold between sets of operators that are not necessarily representations. Such an operator S is called an *intertwining operator*. With this concept we can formulate a theorem that strictly speaking does not deal with representations but with intertwining operators: *Schur's lemma*.

Proposition. Let M and N be two sets of nonsingular linear transformations in spaces V (dimension n) and W (dimension m), respectively. Suppose that both sets are irreducible (the only invariant subspaces are the full space and the origin). Let S be a linear transformation from V to W such that SM = NS. Then either S is the null operator or S is nonsingular and $SMS^{-1} = N$.

Proof: Consider the image of V under S: $Im_S V \subseteq W$. That means that $S\mathbf{r} \in Im_S V$ for all $\mathbf{r} \in V$. This implies that $NS\mathbf{r} = SM\mathbf{r} \in Im_S V$. Therefore, $Im_S V$ is an invariant subspace of W under N. Because N is irreducible, either $Im_S V = 0$ or $Im_S V = W$. In the first case, S is the null operator. In the second case, notice that the kernel of S, the subspace of V mapped on the null vector of W, is an invariant subspace of V under M: if $S\mathbf{r} = 0$ then $NS\mathbf{r} = 0$. Again, because of the irreducibility, either Ker_S is the whole of V, and then S is again the null operator, or $Ker_S = 0$. In the latter case, S is a one-to-one mapping and therefore nonsingular. Therefore, either S is the null operator or it is an isomorphism between the vector spaces V and W, which are then both of dimension n. With respect to bases in the two spaces, the operator S corresponds to a nonsingular matrix and $M = S^{-1}NS$.

This is a very fundamental theorem. Consequences of the theorem are:

(1) If N and M are nonequivalent irreducible representations and SM = NS, then S = 0.

(2) If a matrix S is singular and links two irreducible representations of the same dimension, then S = 0.

(3) A matrix S that commutes with all matrices of an irreducible complex representation is a multiple of the identity. Suppose that an $n \times n$ matrix S commutes with all matrices of a complex irreducible representation. S can be singular and is then the null matrix, or it is nonsingular. In the latter case it has an eigenvalue $\lambda \neq 0$ and $S - \lambda E$ commutes with all the matrices. However, $S - \lambda E$ is singular and therefore the null matrix: $S = \lambda E$. This reasoning is only valid in a complex space, because, generally, the eigenvalues λ are complex.

1.2.2.3. General tensors

Suppose a group K acts linearly on a d-dimensional space V: for any $v \in V$ one has

$$Rv \in V \quad \forall R \in K, v \in V.$$

For a basis \mathbf{a}_i in V this gives a matrix group $\Gamma(K)$ via

$$R\mathbf{a}_i = \sum_{j=1}^d \Gamma(R)_{ji} \mathbf{a}_j, \quad R \in K.$$
(1.2.2.14)

The matrix group $\Gamma(K)$ is a matrix representation of the group *K*. Consider now a linear function *f* on *V*. Because

$$f\left(\sum_{i=1}^{d} \xi_i \mathbf{a}_i\right) = \sum_{i=1}^{d} \xi_i f(\mathbf{a}_i),$$

the function is completely determined by its value on the basis vectors \mathbf{a}_i . A second point is that these linear functions form a vector space because for two functions f_1 and f_2 the function $\alpha_1 f_1 + \alpha_2 f_2$ is a well defined linear function. The vector space is called the *dual space* and is denoted by V^* . A basis for this space is given by functions f_1, \ldots, f_d such that

 $f_i(\mathbf{a}_i) = \delta_{ii},$

because any linear function f can be written as a linear combination of these vectors with as coefficients the value of f on the basis vectors \mathbf{a}_i :

$$f = \sum_{i=1}^{d} f(\mathbf{a}_i) f_i \Leftrightarrow f(\sum_{k=1}^{d} \xi_k \mathbf{a}_k) = \sum_{k=1}^{d} \xi_k \sum_{i=1}^{d} f(\mathbf{a}_i) f_i(\mathbf{a}_k) = \sum_{k=1}^{d} \xi_k f(\mathbf{a}_k).$$

references