tions involving a spontaneous distortion of the crystal lattice that entails a change of shape of the crystallographic or conventional unit cell (Wadhawan, 2000). Such a transformation is accompanied by a change in the number of independent nonzero components of a symmetric second-rank tensor u that describes spontaneous strain.

In discussing ferroelastic and non-ferroelastic domain structures, the concepts of crystal family and holohedry of a point group are useful (*IT* A, 2002). Crystallographic point groups (and space groups as well) can be divided into seven crystal systems and six *crystal families* (see Table 3.4.2.2). A symmetry descent within a crystal family does not entail a qualitative change of the spontaneous strain – the number of independent nonzero tensor components of the strain tensor u remains unchanged.

We shall denote the crystal family of a group M by the symbol Fam$M$. Then a simple criterion for a ferroic phase transition with symmetry descent $G \subset F$ to be a *non-ferroelastic phase transition* is

$$F \subset G, \quad \mathrm{Fam}F = \mathrm{Fam}G. \qquad (3.4.2.25)$$

A necessary and sufficient condition for a *ferroelastic phase transition* is

$$F \subset G, \quad \mathrm{Fam}F \neq \mathrm{Fam}G. \qquad (3.4.2.26)$$

A *ferroelastic domain state* $\mathbf{R}_i$ is defined as a state with a homogeneous spontaneous strain $u^{(i)}$. [We drop the suffix 's' or '(s)' if the serial number of the domain state is given as the superscript (i). The definition of spontaneous strain is given in Section 3.4.3.6.1.] Different ferroelastic domain states differ in spontaneous strain. The symmetry of a ferroelastic domain state $\mathbf{R}_i$ is specified by the stabilizer $I_G(u^{(i)})$ of the spontaneous strain $u^{(i)}$ of the principal domain state $\mathbf{S}_i$ [see (3.4.2.16)]. This stabilizer, which we shall denote by A_i, can be expressed as an intersection of the parent group G and the holohedry of group F_i, which we shall denote HolF_i (see Table 3.4.2.2):

$$A_i \equiv I_G(u^{(i)}) = G \cap \mathrm{Hol}F_i. \qquad (3.4.2.27)$$

This equation indicates that the ferroelastic domain state $\mathbf{R}_i$ has a prominent single-domain orientation. Further on, the term 'ferroelastic domain state' will mean a 'ferroelastic domain state in single-domain orientation'.

In our illustrative example,

$$\begin{aligned} A_1 &= I_{4_z/m_zm_xm_{xy}}(u_{11} - u_{22}) \\ &= \mathrm{Hol}(2_xm_ym_z) \cap m4_z/m_zm_xm_{xy} \\ &= m_xm_ym_z \cap 4_z/m_zm_xm_{xy} = m_xm_ym_z. \end{aligned}$$

The number n_a of ferroelastic domain states is given by

$$n_a = [G : A_1] = |G| : |A_1|. \qquad (3.4.2.28)$$

In our example, $n_a = |4_z/m_zm_xm_{xy}| : |m_xm_ym_z| = 16 : 8 = 2$. In Table 3.4.2.7, last column, the number n_a of ferroelastic domain states is given for all possible ferroic phase transitions.

The number d_a of principal domain states compatible with one ferroelastic domain state (degeneracy of ferroelastic domain states) is given by

$$d_a = [A_1 : F_1] = |A_1| : |F_1|. \qquad (3.4.2.29)$$

In our example, $d_a = |m_xm_ym_z| : |2_xm_ym_z| = 8 : 4 = 2$, *i.e.* two non-ferroelastic principal domain states are compatible with each of the two ferroelastic domain states (*cf.* Fig. 3.4.2.2).

The product of n_a and d_a is equal to the number n of all principal domain states [see equation (3.4.2.19)],

$$n_ad_a = [G : A_1][A_1 : F_1] = [G : F_1] = n. \qquad (3.4.2.30)$$

The number d_a of principal domain states in one ferroelastic domain state can be calculated for all ferroic phase transitions from the ratio of numbers n and n_a that are given in Table 3.4.2.7.

According to Aizu (1969), we can recognize three possible cases:

(i) *Full ferroelastics*: All principal domain states differ in spontaneous strain. In this case, $n_a = n$, *i.e.* $A_1 = F_1$, ferroelastic domain states are identical with principal domain states.

(ii) *Partial ferroelastics*: Some but not all principal domain states differ in spontaneous strain. A necessary and sufficient condition is $1 < n_a < n$, or, equivalently, $F_1 \subset A_1 \subset G$. In this case, ferroelastic domain states are degenerate secondary domain states with degeneracy $n > d_a = |A_1| : |F_1| > 1$. In this case, the phase transition $G \supset F_1$ can also be classified as an *improper ferroelastic* one (see Section 3.1.3.2).

(iii) *Non-ferroelastics*: All principal domain states have the same spontaneous strain. The criterion is $n_a = 1$, *i.e.* $A_1 = G$.

A similar classification for ferroelectric domain states is given below. Both classifications are summarized in Table 3.4.2.3.

Example 3.4.2.1. Domain states in leucite. Leucite ($KAlSi_2O_6$) (see *e.g.* Hatch *et al.*, 1990) undergoes at about 938 K a ferroelastic phase transition from cubic symmetry $G = m\bar{3}m$ to tetragonal symmetry $L = 4/mmm$. This phase can appear in $|G = m\bar{3}m| : |4/mmm| = 3$ single-domain states, which we denote $\mathbf{R}_1$, $\mathbf{R}_2$, $\mathbf{R}_3$. The symmetry group of the first domain state $\mathbf{R}_1$ is $L_1 = 4_x/m_xm_ym_z$. This group equals the stabilizer $I_G(u^{(1)})$ of the spontaneous strain $u^{(1)}$ of $\mathbf{R}_1$ since $\mathrm{Hol}(4_x/m_xm_ym_z) = 4_x/m_xm_ym_z$ (see Table 3.4.2.2), hence this phase is a full ferroelastic one.

At about 903 K, another phase transition reduces the symmetry $4/mmm$ to $F = 4/m$. Let us suppose that this transition has taken place in a domain state $\mathbf{R}_1$ with symmetry $L_1 = 4_x/m_xm_ym_z$; then the room-temperature ferroic phase has symmetry $F_1 = 4_x/m_x$. The $4_x/m_xm_ym_z \supset 4_x/m_x$ phase transition is a non-ferroelastic one $[\mathrm{Hol}(4_x/m_x) = \mathrm{Hol}(4_x/m_xm_ym_z) = 4_x/m_xm_ym_z]$ with $|4_x/m_xm_ym_z| : |4_x/m_x| = 8 : 4 = 2$ non-ferroelastic domain states, which we denote $\mathbf{S}_1$ and $\mathbf{S}_2$. Similar

Table 3.4.2.2. *Crystal systems, holohedries, crystal families and number of spontaneous strain components*

Point group M	Crystal system	Holohedry Hol$M$	Spontaneous strain components: Independent	Spontaneous strain components: Nonzero	Crystal family Fam$M$
23, $m\bar{3}$, 432, $\bar{4}3m$, $m\bar{3}m$	Cubic	$m\bar{3}m$	1	3	Cubic
6, $\bar{6}$, $6/m$, 622, $6mm$, $\bar{6}2m$, $6/mmm$	Hexagonal	$6/mmm$	2	3	Hexagonal
3, $\bar{3}$, 32, $3m$, $\bar{3}m$	Trigonal	$\bar{3}m$	2	3	
4, $\bar{4}$, $4/m$, 422, $4mm$, $\bar{4}2m$, $4/mmm$	Tetragonal	$4/mmm$	2	3	Tetragonal
222, $mm2$, mmm	Orthorhombic	mmm	3	3	Orthorhombic
2, m, $2/m$	Monoclinic	$2/m$	4	4	Monoclinic
1, $\bar{1}$	Triclinic	$\bar{1}$	6	6	Triclinic