

15.1. PHASE IMPROVEMENT BY ITERATIVE DENSITY MODIFICATION

If we use two threshold values, ρ_1 and ρ_2 , to divide the three density regions, the complete formula can be expressed as

$$P(\rho) = \begin{cases} N \exp\left[-(\rho - \bar{\rho})^2/2\sigma^2\right] & \text{for } 2\rho \leq \rho_2 \\ N(a\rho^3 + b\rho^2 + c\rho + d) & \text{for } 2\rho_2 < \rho \leq \rho_1 \\ N(A/\rho)[\ln(\rho_0/\rho)]^{1/2} & \text{for } 2\rho_1 < \rho \leq \rho_0 \end{cases} \quad (15.1.2.11)$$

The parameters a, b, c, d in the cubic polynomial are calculated by matching function values and gradients at ρ_1 and ρ_2 . The parameters in the histogram formula, $\bar{\rho}, \sigma, A, \rho_0, \rho_1, \rho_2$, can be obtained from histograms of known structures.

15.1.2.2.3. The process of histogram matching

Zhang & Main (1990a) demonstrated that, at better than 4 Å resolution, the histogram for an MIR map is generally significantly different from the ideal distribution calculated from atomic coordinates. The obvious course is therefore to alter the map in such a way as to make its density histogram equal to the ideal distribution. Unfortunately, there are an infinite number of maps corresponding to any chosen density distribution, so we must choose a systematic method of altering the map.

The conventional method of performing such a modification is to retain the ordering of the density values in the map. The highest point in the original map will be the highest point in the modified map, the second highest points will correspond in the same way, and so on.

Mathematically, this transformation is represented as follows. Let $P(\rho)$ be the current density histogram and $P'(\rho)$ be the desired distribution, normalized such that their sums are equal to 1. The cumulative distribution functions, $N(\rho)$ and $N'(\rho)$, may then be calculated:

$$N(\rho) = \int_{\rho_{\min}}^{\rho} P(\rho) \, d\rho, \quad (15.1.2.12)$$

$$N'(\rho') = \int_{\rho_{\min}}^{\rho'} P'(\rho) \, d\rho.$$

The cumulative distribution function of a variable transforms a value chosen from the distribution into a number between 0 and 1, representing the position of that value in an ordered list of values chosen from the distribution.

The transformation may, therefore, be performed in two stages. A density value is taken from the initial distribution and the cumulative distribution function of the initial distribution is applied to obtain the position of that value in the distribution. The inverse of the cumulative distribution function for the desired distribution is applied to this value to obtain the density value for the corresponding point in the desired distribution. Thus, given a density value, ρ , from the initial distribution, the modified value, ρ' , is obtained by

$$\rho' = N'^{-1}[N(\rho)]. \quad (15.1.2.13)$$

The distribution of ρ' will then match the desired distribution after the above transformation. The transformation of an electron-density value by this method is illustrated in Fig. 15.1.2.3. The transformation in equation (15.1.2.13) can be achieved through a linear transform represented by

$$\rho'_i = a_i\rho_i + b_i, \quad (15.1.2.14)$$

where $i = \{1, \dots, n\}$ and n is the number of density bins. The above linear transform is sufficient if the number of density bins is large enough. An n value of about 200 is usually quite satisfactory.

Various properties of the electron density are specified in the density histogram, such as the minimum, maximum and mean density, the density variance, and the entropy of the map. The mean density of the ideal map can be obtained by

$$\bar{\rho} = \int_{\rho_{\min}}^{\rho_{\max}} \rho P(\rho) \, d\rho. \quad (15.1.2.15)$$

The variance of the density in the ideal map can be obtained by

$$\sigma(\rho) = \left(\overline{\rho^2} - \bar{\rho}^2\right)^{1/2}, \quad (15.1.2.16)$$

where

$$\overline{\rho^2} = \int_{\rho_{\min}}^{\rho_{\max}} \rho^2 P(\rho) \, d\rho. \quad (15.1.2.17)$$

The entropy of the ideal map can be calculated by

$$S = - \int_{\rho_{\min}}^{\rho_{\max}} P(\rho) \rho \ln(\rho) \, d\rho. \quad (15.1.2.18)$$

Therefore, the process of histogram matching applies a minimum and a maximum value to the electron density, imposes the correct mean and variance, and defines the entropy of the new map. The order of electron-density values remains unchanged after histogram matching.

Histogram matching is complementary to solvent flattening since it is applied to the protein region of a map, whereas

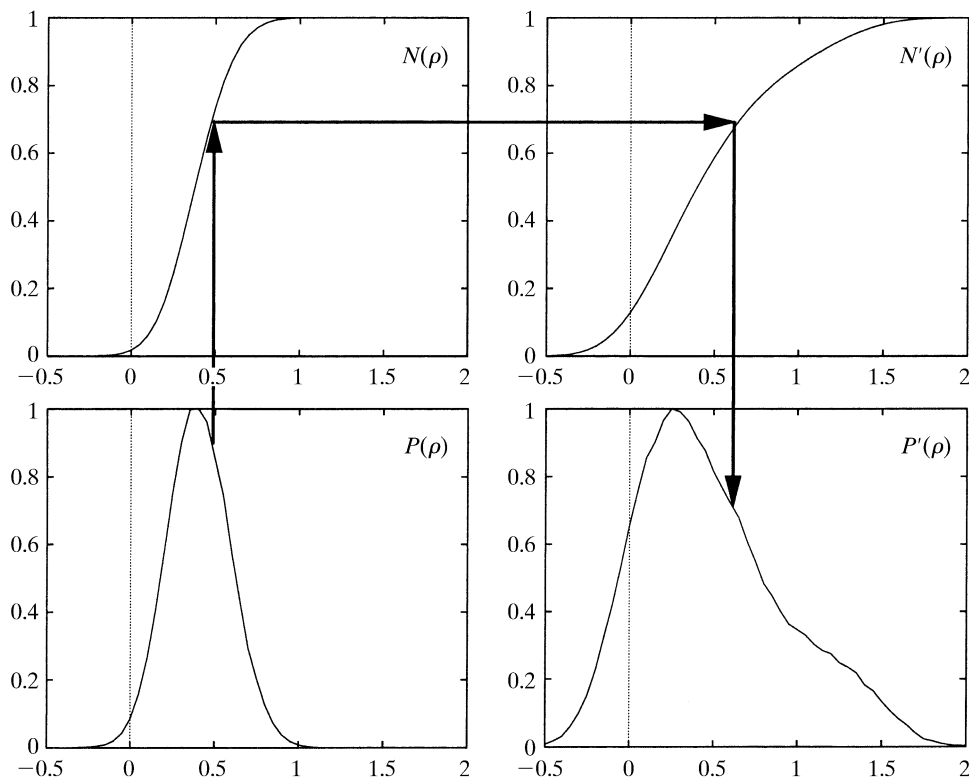


Fig. 15.1.2.3. Transformation of density ρ to ρ'_{mod} by histogram matching.

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solvent flattening only operates on the solvent region of the map. The same envelope that was used for isolating the solvent region can be used to determine the protein region of the cell. An alternative approach is to define separate solvent and protein masks, with uncertain regions excluded from either mask and allowed to keep their unmodified values.

15.1.2.2.4. Scaling the observed structure-factor amplitudes according to the ideal density histogram

In the process of density modification, electron density or structure factors from different sources are compared and combined. It is, therefore, crucial to ensure that all the structure factors and maps are on the same scale. The observed structure factors can be put on the absolute scale by Wilson statistics (Wilson, 1949) using a scale and an overall temperature factor. This is accurate when atomic or near atomic resolution data are available. The scale and overall temperature factor obtained from Wilson statistics are less accurate when only medium- to low-resolution data are available. A more robust method of scaling non-atomic resolution data is through the density histogram (Cowtan & Main, 1993; Zhang, 1993).

The ideal density histogram defines the mean and variance of an electron density, as shown in equations (15.1.2.15) and (15.1.2.16). We can scale the observed structure-factor amplitudes to be consistent with the target histogram using the following formula, obtained from the structure-factor equation and Parseval's theorem. The mean density and the density variance of the observed map can be calculated as

$$\bar{\rho} = (1/V)F(000), \quad (15.1.2.19)$$

$$\sigma^2(\rho) = (1/V) \left[\sum_{\mathbf{h}} |F(\mathbf{h})|^2 \right]^{1/2}. \quad (15.1.2.20)$$

The mean and variance of the electron-density map at the desired resolution are calculated using the target histogram, the mean value of the solvent density, $\bar{\rho}_{\text{solv}}$, and the solvent volume of the cell, V_{solv} . The $F(000)$ term can then be evaluated from equations (15.1.2.15) and (15.1.2.19):

$$F(000) = (V - V_{\text{solv}})\bar{\rho} + V_{\text{solv}}\bar{\rho}_{\text{solv}}. \quad (15.1.2.21)$$

The scale of the observed amplitudes can be obtained from equations (15.1.2.16) and (15.1.2.20),

$$F'(\mathbf{h}) = KF(\mathbf{h}), \quad (15.1.2.22)$$

where

$$K = \left[(\bar{\rho}^2 - \bar{\rho}_{\text{solv}}^2) \right]^{1/2} / \left\{ (1/V) \left[\sum_{\mathbf{h}} |F(\mathbf{h})|^2 \right]^{1/2} \right\}. \quad (15.1.2.23)$$

This method is adequate for scaling observed structure factors at any resolution.

15.1.2.3. Averaging

The averaging method enforces the equivalence of electron-density values between grid points in the map related by noncrystallographic symmetry. The averaging procedure can filter noise, correct systematic error and even determine the phases *ab initio* in favourable cases (Chapman *et al.*, 1992; Tsao *et al.*, 1992).

15.1.2.3.1. Introduction

Noncrystallographic symmetry (NCS) arises in crystals when there are two or more of the same molecules in one asymmetric unit. Such symmetries are local, since they only apply within a sub-region of a single unit cell. A fivefold axis, for example, must be noncrystallographic, since it is not possible to tessellate objects with

fivefold symmetry. Since the symmetry does not map the crystal lattice back onto itself, the individual molecules that are related by the noncrystallographic symmetry will be in different environments; therefore, the symmetry relationships are only approximate.

Noncrystallographic symmetries provide phase information by the following means. Firstly, the related regions of the map may be averaged together, increasing the ratio of signal to noise in the map. Secondly, since the asymmetric unit must be proportionally larger to hold multiple copies of the molecule, the number of independent diffraction amplitudes available at any resolution is also proportionally larger. This redundancy in sampling the molecular transform leads to additional phase information which can be used for phase improvement.

15.1.2.3.2. The determination of noncrystallographic symmetry

The self-rotation symmetry is now routinely solved by the use of a Patterson rotation function (Rossmann & Blow, 1962). The translation symmetry can be determined by a translation function (Crowther & Blow, 1967) when a search model, either an approximate structure of the protein to be determined or the structure of a homologous protein, is available. The searches of the Patterson rotation and translation functions are achieved typically using fast automatic methods, such as *X-PLOR* (Brünger *et al.*, 1987) or *AMoRe* (Navaza, 1994). In cases where no search model is available or the Patterson translation function is unsolvable, either the whole electron-density map, or a region which is expected to contain a molecule, may be rotated using the rotation solution and used as a search model in a phased translation function (Read & Schierbeek, 1988).

Once the averaging operators are determined, the mask can be determined using the local density correlation function as developed by Vellieux *et al.* (1995). This is achieved by a systematic search for extended peaks in the local density correlation, which must be carried out over a volume of several unit cells in order to guarantee finding the whole molecule. The local correlation function distinguishes those volumes of crystal space which map onto similar density under transformation by the averaging operator. Thus, in the case of improper NCS, a local correlation mask will cover only one monomer. In the case of a proper symmetry, a local correlation mask will cover the whole complex (Fig. 15.1.2.4a,b).

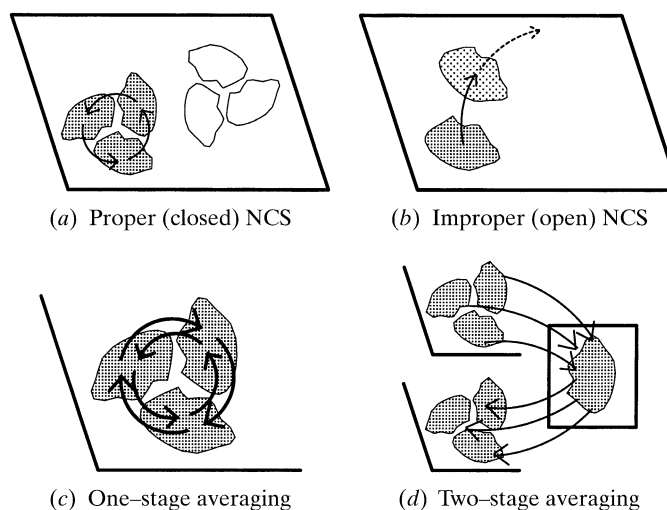


Fig. 15.1.2.4. Types of noncrystallographic symmetry and averaging calculation.